

(Semi-)Analytical solutions of 1D partial differential equations

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Thanks: S. Balbus, J. Papaloizou, D. Lynden-Bell

Outline

- 1 Finding analytical solutions
 - Introduction
 - Fourier
 - JWKB
 - Non-linear waves
 - Self-similar solutions
- 2 Using analytical solutions
 - Test codes
 - New physics
 - Numerical algorithms

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Motivation

Astrophysical gases are a complicated business...

- Rich microphysics
- Radiation transport (cooling, thermal diffusion)
- Chemistry
- Magnetic fields
- High Mach flows
- High resolution features

⇒ Need complicated benchmarks to test codes

The general method

General method

- Consider a set of partial differential equations

$$\mathcal{F}[x, t, y(x, t), \partial_t y, \partial_x y, \partial_{xx} y, \dots] = 0$$

- Choose a form for the solution $y(x, t)$
- Plug it in the equations...

And if it doesn't work ?

- Adjust parameters of the solution
- Choose better initial or boundary conditions
- Fudge the physics (change \mathcal{F})

Fourier modes

A Fourier mode

$$y = y_0 e^{(s.t + ik.x)}$$

- A must for linear equations...

Example

The heat transport equation

$$\partial_t y = \partial_{xx} y \Rightarrow s + k^2 = 0$$

(dispersion relation \equiv fitting parameters)

- Can work for non-linear equations as well !

Fourier modes

For non-linear PDEs

Examples

- Torsional Alfvén waves for ideal MHD
- Incompressible MRI modes are solutions of the non-linear isothermal ideal MHD Hill system (Note: it needs initially homogeneous conditions).
- For linear cooling and a proper choice of parameters, you get solutions with cooling, resistivity and viscosity as well (Lesaffre & Balbus, 2008) !

JWKB approximation

JWKB expansion

$$y = \exp\left[st + \frac{1}{\varepsilon} \sum_{n=0}^{n=\infty} \varepsilon^n S_n(x)\right]$$

- Find constraining ODEs for the S_n at each order in ε .
- Then let $\varepsilon \rightarrow 1$.
- The series is sometimes finite or converges very fast.

Example

Sheared waves for the hydrodynamical Hill system (Balbus & Hawley 2006, ApJ)

Non-Linear waves

Non-linear wave form

$$y = Y(s.t - k.x)$$

Injecting this solution can convert the PDE into an ODE for Y.

Examples

- HD sheared waves (Fromang, Papaloizou 2007, A&A)
- Analytical Steady-state shocks (Lesaffre 2006, GAFFD)
- Method of characteristics

Self-Similar solutions.

General self-similar form

$$y = f(t) + g(t).Y[a(t) + b(t).x]$$

- Inject that form in the PDE
- Then request variable separation between t and $X = a(t) + b(t).x$
- That yields one ODE for each $a(t)$, $b(t)$, $f(t)$, $g(t)$ and $Y(X)$

Notes

- Non-linear waves are obtained by this procedure
- Variable separation is more general than dimensional analysis
- Although the time coefficients are often found to be power laws or exponentials

Detailed example

Adimensional equation for convective/radiative transport

$$\partial_t \ln T = \partial_X \mathcal{L}(\partial_X \ln T) \text{ where } \mathcal{L}(\nabla) = \nabla + U \cdot [\nabla - \nabla_a]_+^{3/2}$$

- We first inject the form $\ln T = f(t)F[a(t)x]$ into the equation:

$$\dot{f} \left(F + \frac{f}{\dot{f}} \frac{\dot{a}}{a} X \partial_X F \right) = a \partial_X \mathcal{L}(f a \partial_X F) \text{ where } X = ax$$

- Request variable separation \Rightarrow get time constraints

$$\begin{cases} f \dot{f} = \alpha \\ f a = \beta \end{cases} \text{ with } \alpha \text{ and } \beta \text{ constants}$$

- and a second order ODE for the shape function $F(X)$:

$$\alpha(F - X \partial_X F) = \beta \partial_X \mathcal{L}(\beta \partial_X F)$$

Notes

- You get the solution for any function \mathcal{L} (pick up your favourite convection theory)
- $a(t) \propto t^{-2}$ at large times: recover dimensional analysis result for the radiative case
- I believe the present family of solutions has dimension $(3+2-2)$
- You can find broader families by using all $a(t)$, $b(t)$, $f(t)$, $g(t)$
- You can include nuclear heating provided it is of the form $\varepsilon = a(t)E(X)$

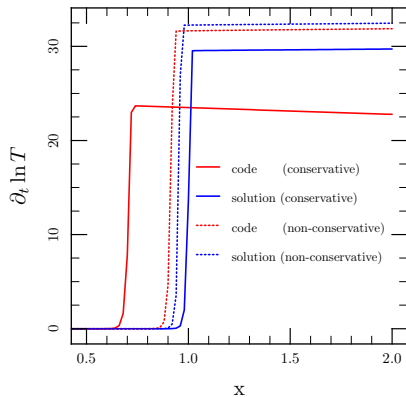
Exercise

Apply technique to Navier-Stokes equations and recover shock-tube, blast waves, bubbles, exponential solutions (c.f. Zel'dovich & Raizer) with viscosity.

Code Results

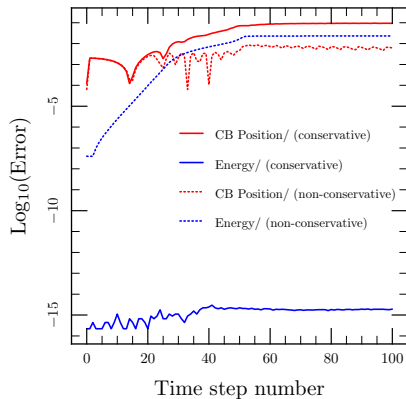
Comparing two ways of discretizing the energy equation

conservative vs. non-conservative



(a) Snapshot after 80 time steps

conservative vs. non-conservative



(b) Evolution of errors

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Testing codes

Trust a code's result only when you already know it...

- Check codes, debug them
- Improve their accuracy
- Find their domains of application
- Probe their convergence properties

Example

Lesaffre & Balbus (2007) probe the dissipation properties of ZEUS thanks to Alfvén waves with viscosity and resistivity. They measure a scaling law for the total dissipation:

$$\eta_N + \nu_N = 0.76 \left(\frac{k}{2\pi}\right)^{1.6} \Delta x^2 \beta^{-1/2} + 1.08 \Delta x C \beta^{-1}$$

where k is the wave number, Δx is the space resolution, C is the Courant number and β is the plasma- β parameter

Adress new physics

- Basis for linear analysis
 - parasitic modes
 - frozen modes
- Get physical behaviour for a wide range of parameters
- Understand continuity / discontinuity of solutions

Example

$\partial_{xx} T$ and $\partial_x N$ are discontinuous at convective/radiative boundaries (on self-similar solutions of the heat and chemical transport problem)

Build new numerical algorithms

- First find a large family of solutions to your problem
- Carefully assess the dimension of that class (beware of degeneracy, numerical and mathematical)
- The analytical solutions space needs to be dense in the solutions space
- Carefully design a way of picking up only one analytical solution (given initial and boundary conditions)
- Assess the stability of the process

Example: Godunov schemes

- Assume initial conditions=smooth+Riemann problem.
- Evolve thanks to self-similar solution.
- Reconstruct the profile to get smooth+Riemann

Summary

- It is possible to find **large classes** of (semi-)analytical solutions, even for complicated problems.
- They can be **really useful**...
 - to probe codes
 - to understand physics
 - to build new algorithms